

# Relativistic nature of the Lorentz force.

(1)

S - frame

$$F = \frac{1}{4\pi\epsilon_0 c^2} q \frac{2I}{r} \delta, \quad \text{Current } I = \int j ds = j A = \rho_- \delta A$$

↑  
cross section

$$[S] F = \frac{1}{4\pi\epsilon_0 c^2} \frac{2q\delta}{r} \cdot \rho_- \delta A = \frac{1}{4\pi\epsilon_0 c^2} \frac{\delta^2}{r} \frac{2q\rho_- A}{r}$$

S' - frame

Charge density depends on the volume which changes due to the relativistic contraction, i.e.

$$\underbrace{\rho_0 L_0 A}_{(\text{at rest}) Q} = \underbrace{\rho L_0 \sqrt{1 - \frac{\delta^2}{c^2}} A}_{(\delta \neq 0) Q} \Rightarrow \rho = \frac{\rho_0}{\sqrt{1 - \frac{\delta^2}{c^2}}}$$

(since charge is invariant)

$\rho_+$  is at rest in the frame S  $\equiv \rho_0$

$$(1) \rho_+' = \frac{\rho_+}{\sqrt{1 - \frac{\delta^2}{c^2}}}; \quad \rho_- \text{ is at rest in frame } S', \text{ i.e. } \rho_-' \equiv \rho_0$$

$$(2) \rho_-' = \rho_- \sqrt{1 - \frac{\delta^2}{c^2}}$$

$$\rho_- = \rho_+$$

$$\rho' = \rho_+' + \rho_-' = \rho_+ \left( \frac{1}{\sqrt{1 - \frac{\delta^2}{c^2}}} - \sqrt{1 - \frac{\delta^2}{c^2}} \right) = \rho_+ \frac{\frac{\delta^2}{c^2}}{\sqrt{1 - \frac{\delta^2}{c^2}}}$$

$$E' = \frac{\rho' A}{\pi \epsilon_0 r} \quad (\text{see Blatt 11})$$

(2)

$$= \frac{\rho A}{\pi \epsilon_0 r} \frac{\frac{\delta^2}{c^2}}{\sqrt{1 - \frac{\delta^2}{c^2}}}$$

$$[S'] F' = q E' = \frac{q}{\pi \epsilon_0} \frac{\rho A}{r} \frac{\frac{\delta^2}{c^2}}{\sqrt{1 - \frac{\delta^2}{c^2}}}$$

$F' = F$  at least for small velocities

if velocities are high then:

$$\left\{ \begin{array}{l} \Delta p' = F' \Delta t' \\ \Delta p = F \Delta t \end{array} \right. \quad \text{and } \Delta p' = \Delta p \text{ independent of the frame}$$

$$F' \Delta t' = F \Delta t$$

However due to the time contraction

$$\Delta t' = \Delta t \sqrt{1 - \frac{v^2}{c^2}}$$

Thus  $F' \Delta t' = F \Delta t$  as expected

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Lorentz - Transformation des elektromagnetischen Feldes (3)

$$\begin{cases} ct' = \gamma ct - \beta x \\ x' = -\beta ct + \gamma x \\ y' = y \\ z' = z \end{cases}$$

$$\beta = \frac{v}{c}$$

$$\gamma = \frac{1}{\sqrt{1 - \beta^2}}$$

a)  $\begin{pmatrix} ct' \\ x' \\ y' \\ z' \end{pmatrix} = \begin{pmatrix} \gamma & -\beta \\ -\beta\gamma & \gamma \\ 0 & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} ct \\ x \\ y \\ z \end{pmatrix}$

$$A_{\mu}^{\nu'} = \begin{pmatrix} \gamma & -\beta & 0 & 0 \\ -\beta\gamma & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} = \begin{pmatrix} L & 0 \\ 0 & E \end{pmatrix}$$

b)  $F^{\mu\nu} = \begin{pmatrix} 0 & -E_x & -E_y & -E_z \\ E_x & 0 & -B_z & B_y \\ E_y & B_z & 0 & -B_x \\ E_z & -B_y & B_x & 0 \end{pmatrix} = \begin{pmatrix} A & B \\ -B^T & C \end{pmatrix}$

$$F^{\mu'\nu'} = A_{\mu}^{\nu'} F^{\mu\nu} (A^T)^{\nu}{}_{\mu'} = \begin{pmatrix} L & 0 \\ 0 & E \end{pmatrix} \begin{pmatrix} A & B \\ -B^T & C \end{pmatrix} \begin{pmatrix} L^T & 0 \\ 0 & E \end{pmatrix} = \begin{pmatrix} L & 0 \\ 0 & E \end{pmatrix} \begin{pmatrix} A L^T & B E \\ -B^T L^T & C E \end{pmatrix}$$

$$= \begin{pmatrix} L A L^T & L B E \\ -E B^T L^T & E C E \end{pmatrix}$$

$$L A L^T = \begin{pmatrix} \gamma & -\beta\gamma \\ -\beta\gamma & \gamma \end{pmatrix} \begin{pmatrix} 0 & -E_x \\ E_x & 0 \end{pmatrix} \begin{pmatrix} \gamma & -\beta\gamma \\ -\beta\gamma & \gamma \end{pmatrix} = \begin{pmatrix} \gamma & -\beta\gamma \\ -\beta\gamma & \gamma \end{pmatrix} \begin{pmatrix} E_x \beta & -E_x \\ E_x & \beta E_x \end{pmatrix} = \begin{pmatrix} E_x \gamma^2 \beta - \beta \gamma^2 E_x & -E_x \gamma^2 + \gamma^2 \beta^2 E_x \\ -E_x \gamma^2 \beta^2 + E_x \gamma^2 & E_x \gamma^2 \beta - \beta \gamma^2 E_x \end{pmatrix} = E_x \gamma^2 \begin{pmatrix} 0 & \beta^2 - 1 \\ -\beta^2 + 1 & 0 \end{pmatrix} = E_x \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$$

$$LB = \begin{pmatrix} 1 & -j\beta \\ -j\beta & 1 \end{pmatrix} \begin{pmatrix} -E_y & -E_z \\ -B_x & B_y \end{pmatrix} = \begin{pmatrix} -E_y + \beta B_x & -E_z - \beta B_y \\ \beta E_y - B_x & \beta E_z + B_y \end{pmatrix} \quad (4)$$

$$F^{\mu\nu'} = \begin{pmatrix} 0 & -E_x & -jE_y + j\beta B_x & -jE_z - j\beta B_y \\ E_x & 0 & j\beta E_y - j\beta B_x & j\beta E_z + j\beta B_y \\ jE_y - j\beta B_x & j\beta E_y + j\beta B_x & 0 & -B_x \\ j\beta E_z + j\beta B_y & -j\beta E_z - j\beta B_y & B_x & 0 \end{pmatrix}$$

Thus,

$$E_x' = E_x$$

$$E_y' = jE_y - j\beta B_x$$

$$E_z' = jE_z + j\beta B_y$$

$$B_x' = B_x$$

$$B_y' = j\beta E_z + j\beta B_y$$

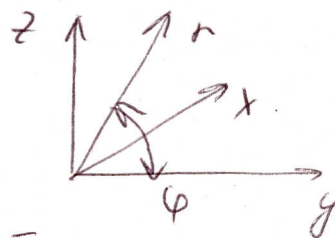
$$B_z' = j\beta E_y - j\beta B_x$$

All the calculations are in CGS!

B Draht mit Ladungsdichte  $\lambda$ :

(3)

a)  $E_r = \frac{2\lambda}{r}$  ;  $E_r = \sqrt{E_y^2 + E_z^2}$



in the rest frame

$$E_y = \frac{2\lambda}{r} \cos\varphi, \quad B_{x,y,z} = 0, \quad E_x = 0$$

$$E_z = \frac{2\lambda}{r} \sin\varphi$$

let the frame move along x. Then:

$$E_x' = E_x, \quad E_y' = \gamma E_y = \gamma \frac{2\lambda}{r} \cos\varphi, \quad E_z' = \gamma E_z = \gamma \frac{2\lambda}{r} \sin\varphi$$

$$B_x' = 0; \quad B_y' = \gamma\beta E_z = \gamma\beta \frac{2\lambda}{r} \sin\varphi, \quad B_z' = -\gamma\beta E_y = -\gamma\beta \frac{2\lambda}{r} \cos\varphi$$

Thus:

$$\vec{E}' = \gamma \frac{2\lambda}{r} (\vec{e}_y \cos\varphi + \vec{e}_z \sin\varphi) = \gamma \frac{2\lambda}{r} \cdot \vec{e}_r = \gamma \vec{E}$$

$$\vec{B}' = \gamma\beta \frac{2\lambda}{r} (\vec{e}_y \sin\varphi - \vec{e}_z \cos\varphi) = -\gamma\beta \frac{2\lambda}{r} \vec{e}_\varphi$$

i.e.  $|\vec{B}'| = B_\varphi = \gamma\beta \frac{2\lambda}{r}$   
 ↑  
 tangential

b)  $v \ll c$

$$\vec{B}' = \vec{B}$$

$$B_\varphi = -\frac{v}{c} \frac{2\lambda}{r} = -\frac{1}{c} \frac{2j}{r}$$